



Détection de Rayonnement à Très Basse Température

Perspectives des détecteurs cryogéniques

Détection des événements rares

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Outline

➤ The quest of neutrino mass and nature

➤ Neutrinoless double beta decay

↳ Cuoricino, CUORE experiments

➤ Direct measurement through single beta decay

↳ MARE experiment

➤ Conclusions and prospects

Crucial role of low-temperature calorimeters

General overview: main pieces of the puzzle

Oscillation experiments



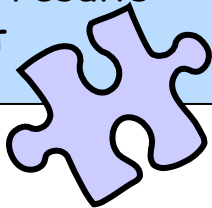
Neutrinos have non-zero mass

ν mass scale is crucial over 2 fronts:
elementary particles physics & cosmology / astrophysics

Cosmology

$$\Sigma m_{\nu i} < 0.7 - 1 \text{ eV}/c^2$$

- very sensitive
- spread in recent results
- model dependent



WE NEED TO KNOW:

- absolute ν mass scale
- nature of ν mass

Direct search through β decay

Potential sensitivity
 $m(\nu_e) < 2.2 \text{ eV}/c^2$

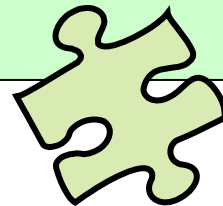
- completely model free
- future sensitivity: 0.2 eV
- in the game if m_ν quasidegenerate



$0\nu\beta\beta$ decay

$$\langle m_\nu \rangle < 0.5 \text{ eV}/c^2$$

- future sensitivity 0.05 eV
- controversial claim: 0.4 eV
- works only if neutrino is a Majorana particle



Neutrino flavor oscillations and mass scale

what **we presently know** from neutrino flavor oscillations

oscillations **do** occur

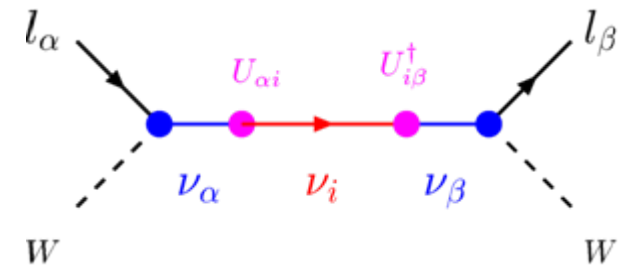


neutrinos are **massive**

Mass eigenstates (ν_1, ν_2, ν_3)

are different from

Flavor eigenstates $(\nu_e, \nu_\mu, \nu_\tau)$



$$\nu_\alpha = \sum_j U_{\alpha j} \nu_j$$

given the three **ν mass eigenvalues** M_1, M_2, M_3 we have

approximate measurements of **two** ΔM_{ij}^2 $(\Delta M_{ij}^2 \equiv M_i^2 - M_j^2)$

$\Delta M_{12}^2 \sim (9 \text{ meV})^2$ **Solar**

$|\Delta M_{23}^2| \sim (50 \text{ meV})^2$ **Atmospheric**

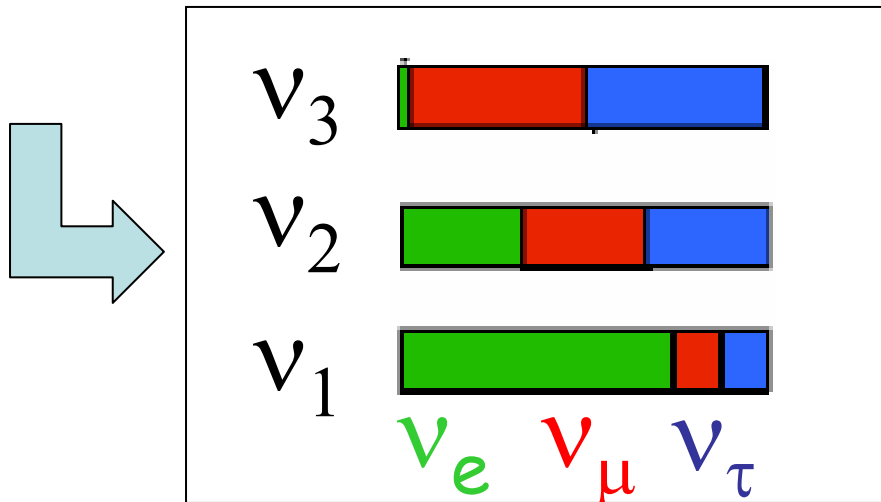
Neutrino mixing

we have also approximate measurements and/or constraints on U_{ij} → elements of the ν mixing matrix

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

parametrized with $c_{ij} \equiv \cos\Theta_{ij}$
 three angles
 and three phases $s_{ij} \equiv \sin\Theta_{ij}$

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13}e^{i\delta} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13}e^{i\delta} \end{pmatrix} \begin{pmatrix} e^{i\alpha_1/2} \nu_1 \\ e^{i\alpha_2/2} \nu_2 \\ \nu_3 \end{pmatrix}$$

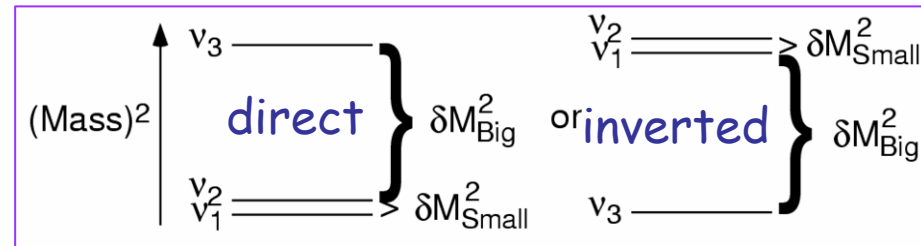


The present knowledge of the values for $U_{\alpha i}$ fixes the composition of the mass eigenstates in terms of flavor eigenstates and vice-versa.

Neutrino flavor oscillations and mass scale

what we do not know from neutrino flavor oscillations:

① neutrino mass **hierarchy** →



② **absolute neutrino mass scale** → **degeneracy?** ($M_1 \sim M_2 \sim M_3$)

③ **DIRAC** or **MAJORANA** nature of neutrinos

Like electrons, neutrinos and antineutrinos are **distinct particles**. Lepton number plays the role of the electric charge.

Neutrinos and antineutrinos are the **same particle**. Their behavior is ruled by helicity.

Tools for the investigation of the ν mass scale

Tools	Present sensitivity	Future sensitivity (a few year scale)
Cosmology (CMB + LSS)	0.7 - 1 eV	0.1 eV
Neutrinoless Double Beta Decay	0.5 eV	0.05 eV
Single Beta Decay	2.2 eV	0.2 eV

Model dependent
Laboratory measurements
(LT calorimeters are key instruments)

Complementarity of cosmology, single and double β decay

Cosmology, **single** and **double β** decay measure different combinations of the neutrino mass eigenvalues, constraining the **neutrino mass scale**

In a standard three active neutrino scenario:

$$\Sigma \equiv \sum_{i=1}^3 M_i$$

cosmology
simple sum
pure kinematical effect

$$\langle M_\beta \rangle \equiv \left(\sum_{i=1}^3 M_i^2 |U_{ei}|^2 \right)^{1/2}$$

beta decay
incoherent sum
real neutrino

$$\langle M_{\beta\beta} \rangle \equiv \left| \sum_{i=1}^3 M_i |U_{ei}|^2 e^{i\alpha_i} \right|$$

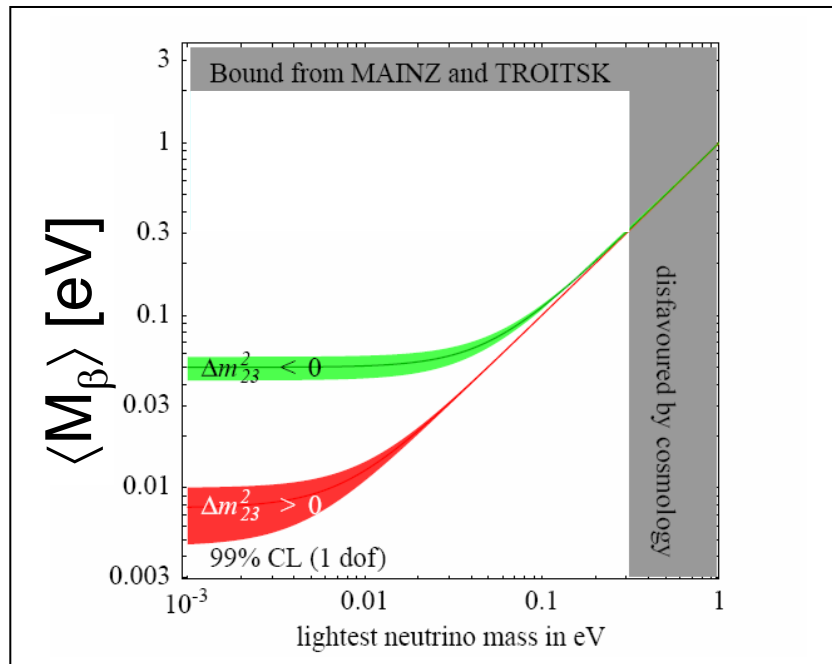
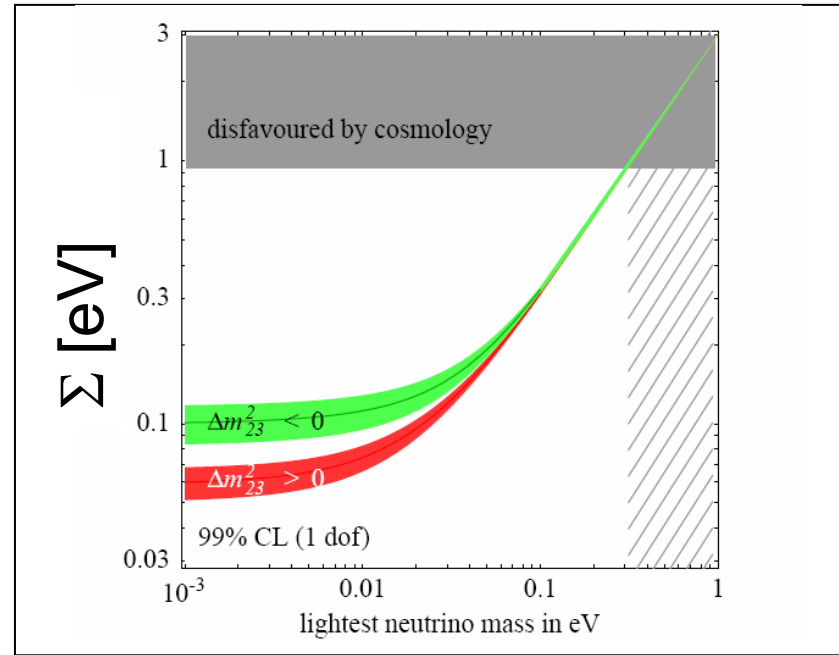
double beta decay
coherent sum
virtual neutrino
Majorana phases

Present bounds

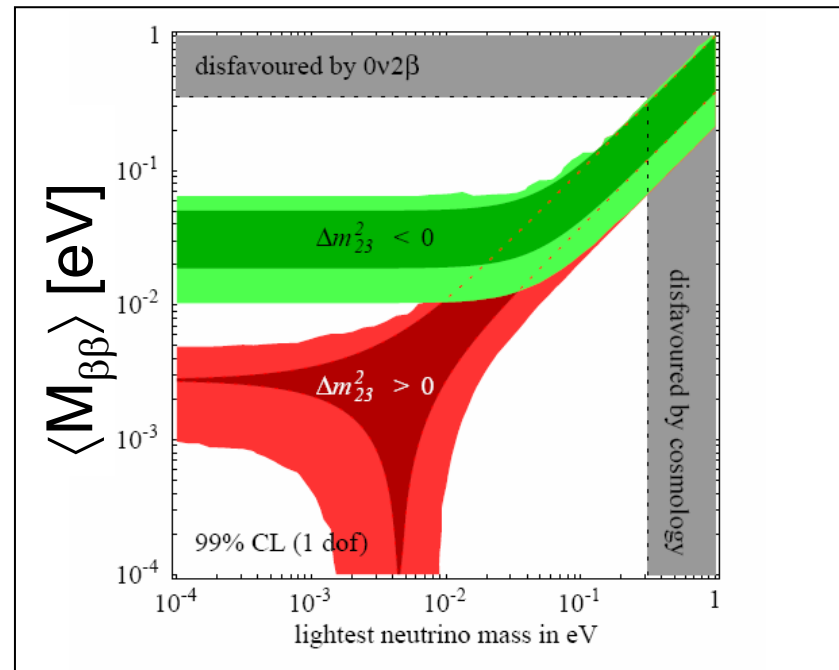
The three constrained parameters can be plotted as a function of the **lightest neutrino mass**

Two bands appear in each plot, corresponding to **inverted** and **direct** hierarchy

The two bands merge in the **degenerate** case (the only one presently probed)



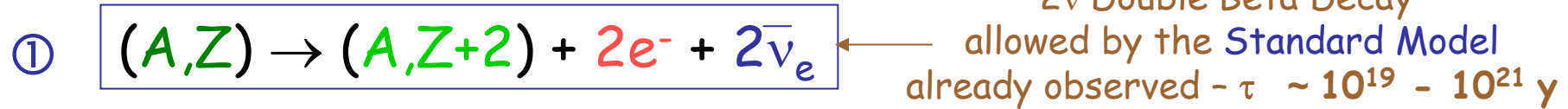
Strumia-Vissani hep-ph/0503246



DOUBLE BETA DECAY

Decay modes for Double Beta Decay

Two decay modes are usually discussed:



Process ② would imply **new physics** beyond the Standard Model

violation of **lepton number conservation**

It is a very **sensitive test to new physics** since the phase space term is much larger for the neutrinoless process than for the standard one

interest for 0ν -DBD lasts for ~ 70 years !

Goepfert-Meyer proposed the standard process in 1935

Racah proposed the neutrinoless process in 1937

0ν -DBD, neutrino mass and neutrino flavor oscillations

Observation of 0ν -DBD



$$m_\nu \neq 0$$
$$\overline{\nu} \equiv \nu$$

how 0ν -DBD is connected to neutrino mixing matrix and masses

neutrinoless
Double Beta Decay
rate

Phase
space

Nuclear
matrix elements

Effective
Majorana mass

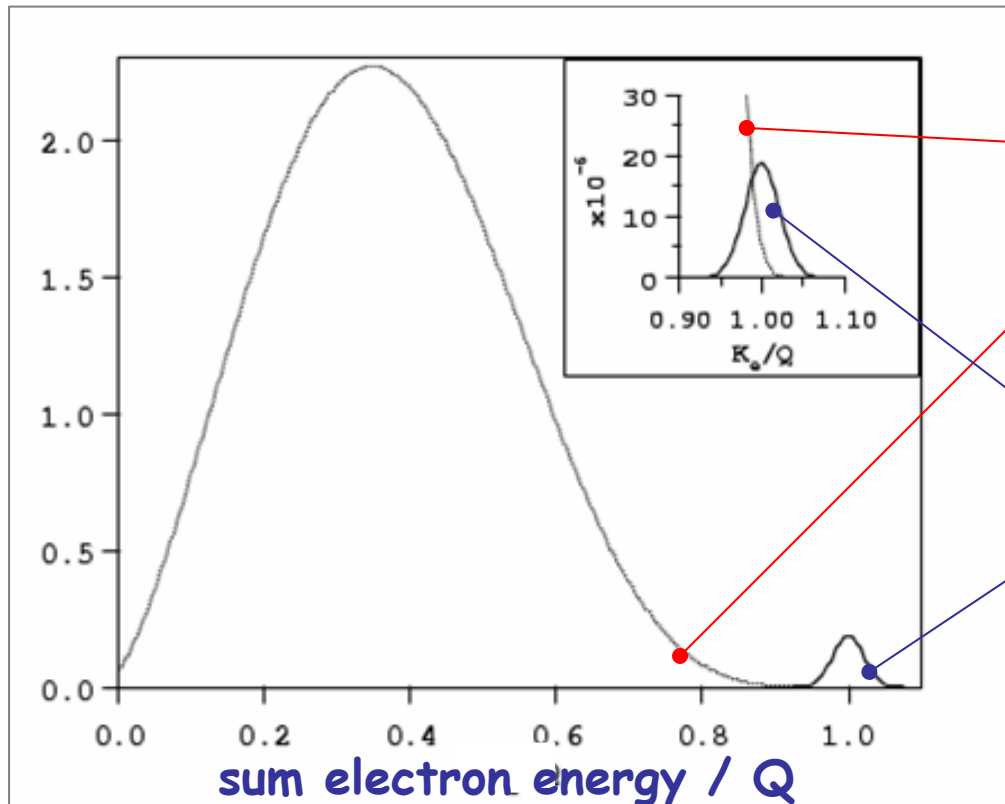
$$1/\tau = G(Q,Z) |M_{\text{nucl}}|^2 \langle M_{\beta\beta} \rangle^2$$

$$\langle M_{\beta\beta} \rangle = \left| |U_{e1}|^2 M_1 + e^{i\alpha_1} |U_{e2}|^2 M_2 + e^{i\alpha_2} |U_{e3}|^2 M_3 \right|$$

can be of the order of $\sim 50 \text{ meV}$ in case of inverted hierarchy

Electron sum energy spectra in DBD

The **shape** of the **two electron sum energy spectrum** enables to distinguish among the two different discussed decay modes



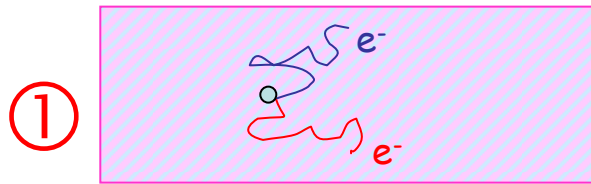
two neutrino DBD
continuum with maximum at $\sim 1/3 Q$

neutrinoless DBD
peak enlarged only by
the detector energy resolution
Typical energy for most
interesting candidates: **2 - 3 MeV**

Experimental approaches to direct searches

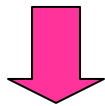
Two approaches for the detection of the two electrons:

Calorimetric



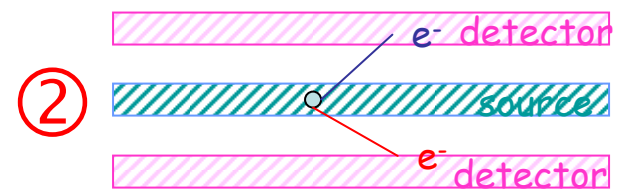
Source \equiv Detector
(calorimetric technique)

- scintillation
- **cryogenic macrocalorimeters (bolometers)**
- solid-state devices
- gaseous detectors



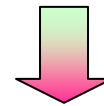
high efficiency and energy resolution

"Tracko-calor"



Source \neq Detector

- scintillation
- gaseous TPC
- gaseous drift chamber
- magnetic field and TOF



event reconstruction,
but low energy resolution

Experimental sensitivity to 0ν -DBD

sensitivity F: lifetime corresponding to the minimum detectable number of events over background at a given confidence level

$b \neq 0$

b: specific background coefficient
[counts/(keV kg y)]

$b = 0$

live time
source mass
energy resolution

$$F \propto (MT / b\Delta E)^{1/2}$$

$$F \propto MT$$

background level

importance of the **nuclide choice**
(but **large uncertainty** due to nuclear physics)

$$\text{sensitivity to } \langle M_{\beta\beta} \rangle \propto (F/Q |M_{\text{nucl}}|^2)^{1/2} \propto \frac{1}{Q^{1/2} |M_{\text{nucl}}|} \left(\frac{b\Delta E}{MT} \right)^{1/4}$$

Present experimental situation in the search for 0ν -DBD

The three presently most sensitive experiments:

- Heidelberg - Moscow (HM) (LNGS) (Source \equiv Detector)
the most sensitive DBD experiment since 10 years (stopped in May 03)
classical technique of high resolution Ge diodes
A subset of the collaboration claims for discovery
- CUORICINO (LNGS) (Source \equiv Detector)
it is an intermediate generation experiment with a sensitivity to neutrino mass similar to HM (stopped in June 08)
innovative technique of low-temperature calorimetry
- NEMO3 (LSM) (Source \neq Detector)
it is an intermediate generation experiment capable to study different candidate nuclides and to improve the HM results (running)
more developed example of "tracko-calo" approach



Neutrinoless DBD

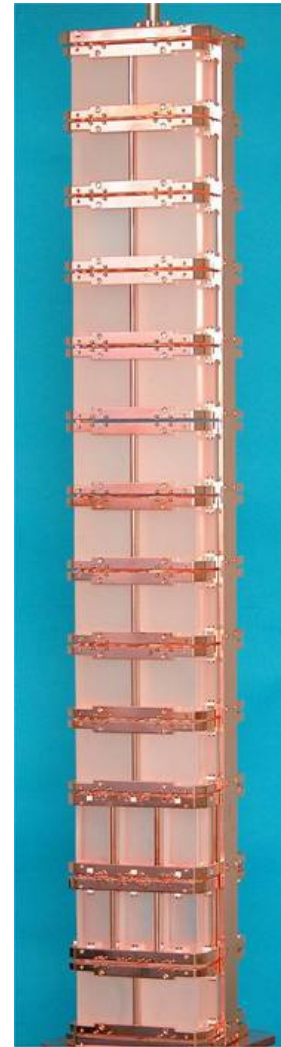
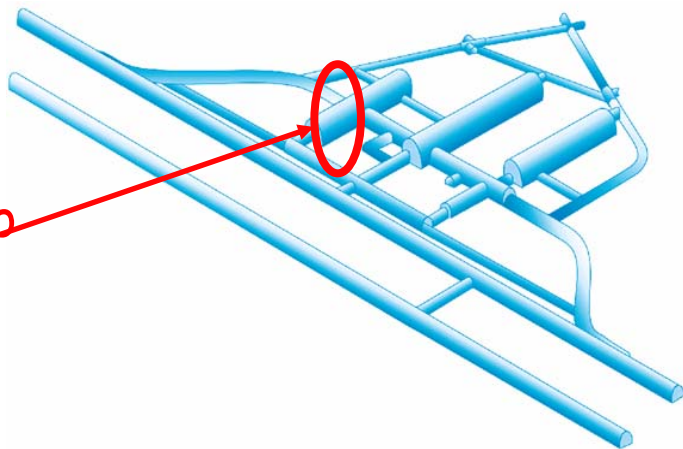
Underground National Laboratory of
Gran Sasso

located in the highway tunnel

3500 m.w.e.

$24\mu / \text{m}^2 / \text{d}$

ITALY



T= 10 mK

total active mass

➤ TeO_2 : 40.7 kg

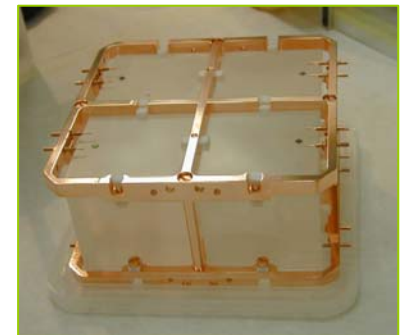
➤ ^{130}Te : 11.3 kg

➤ ^{128}Te : 10.5 kg

- 11 modules with 4 *big* detectors
 - ▼ 44 TeO_2 crystals
 - ▼ $5 \times 5 \times 5 \text{ cm}^3 \Rightarrow 790 \text{ g}$
 - ▷ TeO_2 mass $\Rightarrow 34.76 \text{ kg}$

- 2 modules with 9 *small* detectors
 - ▼ 18 TeO_2 crystals
 - ▼ $3 \times 3 \times 6 \text{ cm}^3 \Rightarrow 330 \text{ g}$
 - ▷ TeO_2 mass $\Rightarrow 5.94 \text{ kg}$
 - 4 crystals are enriched
 - ▼ $2 \times ^{130}\text{TeO}_2 + 2 \times ^{128}\text{TeO}_2$

◆ Heat @ 10 mK with
Ge/NTD thermometer



Scheme of a TeO₂ macrocalorimeter

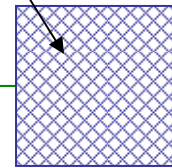
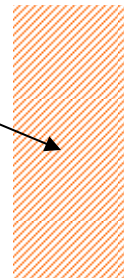
Te dominates in mass the compound
Excellent mechanical and thermal properties

Energy absorber

TeO₂ crystal

$$C \cong 2 \text{ nJ/K} \cong 1 \text{ MeV} / 0.1 \text{ mK}$$

Heat sink
 $T \cong 10 \text{ mK}$



Thermometer
NTD Ge-thermistor
 $R \cong 100 \text{ MW}$
 $dR/dT \cong 100 \text{ kW/mK}$

Thermal coupling

$$G \cong 4 \text{ nW / K} = 4 \text{ pW / mK}$$

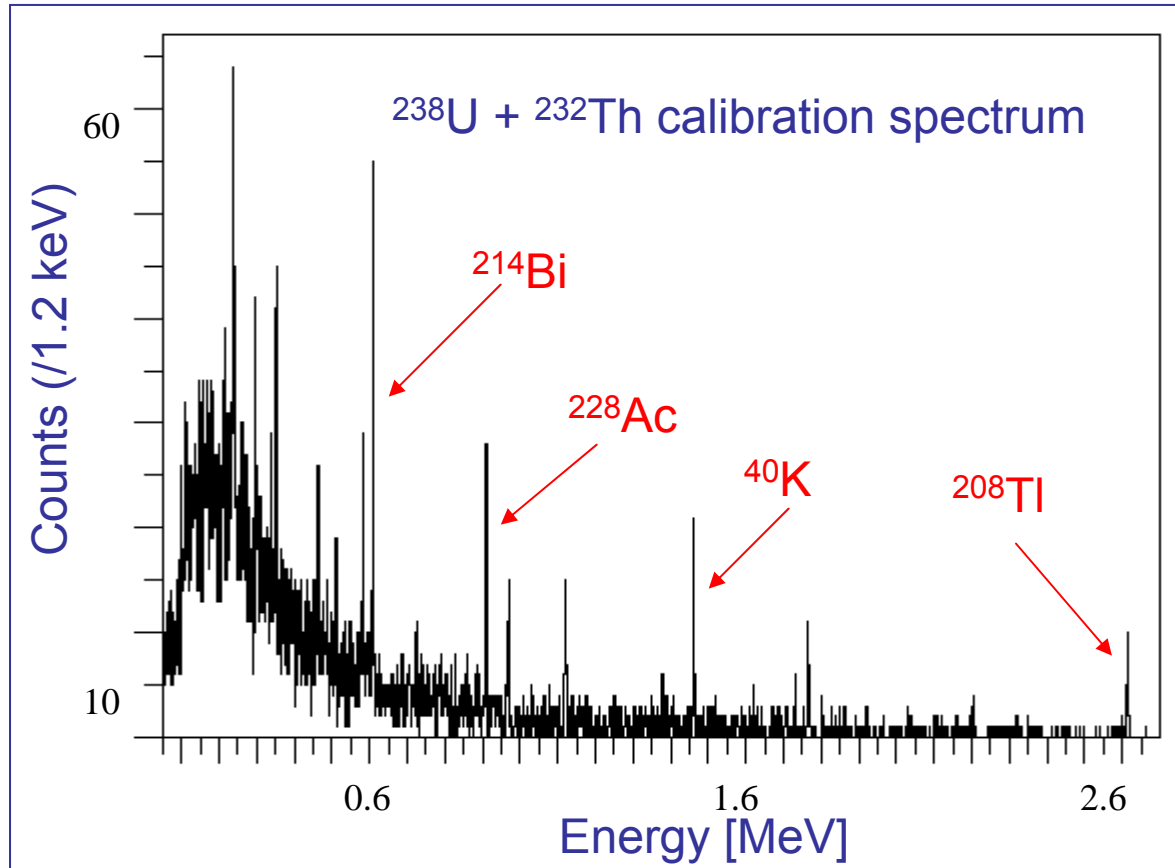
- ◆ Temperature signal: $\Delta T = E/C \cong 0.1 \text{ mK}$ for $E = 1 \text{ MeV}$
- ◆ Bias: $I \cong 0.1 \text{ nA} \Rightarrow$ Joule power $\cong 1 \text{ pW} \Rightarrow$ Temperature rise $\cong 0.25 \text{ mK}$
- ◆ Voltage signal: $\Delta V = I \times dR/dT \times \Delta T \Rightarrow \Delta V = 1 \text{ mV}$ for $E = 1 \text{ MeV}$
- ◆ Signal recovery time: $\tau = C/G \cong 0.5 \text{ s}$
- ◆ Noise over signal bandwidth (a few Hz): $V_{\text{rms}} = 0.2 \text{ mV}$ In real life signal about a factor 2 - 3 smaller

Energy resolution (FWHM): $\cong 1 \text{ keV}$

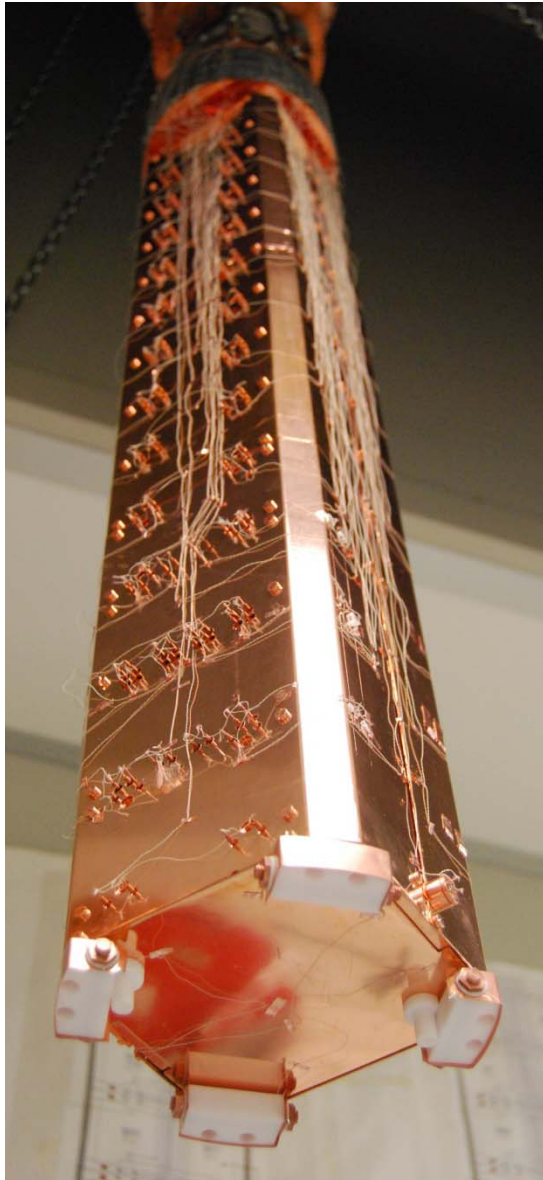
Technical results on detector performances

Performance of CUORICINO-type detectors ($5 \times 5 \times 5 \text{ cm}^3$ - 790 g):

- ◆ Detector base temperature: $\sim 7 \text{ mK}$
- ◆ Detector operation temperature: $\sim 9 \text{ mK}$
- ◆ Detector response: $\sim 250 \text{ mV/ MeV}$
- ◆ FWHM resolution: $\sim 3 \text{ keV @ } 2.6 \text{ MeV}$



In Cuoricini sempiterna memoria



End of June 2008:
Cuoricino has been shut down

Saturation of sensitivity
Need of experimental space in
hallA for further tests



Neutrinoless DBD results

UPDATED RESULTS:

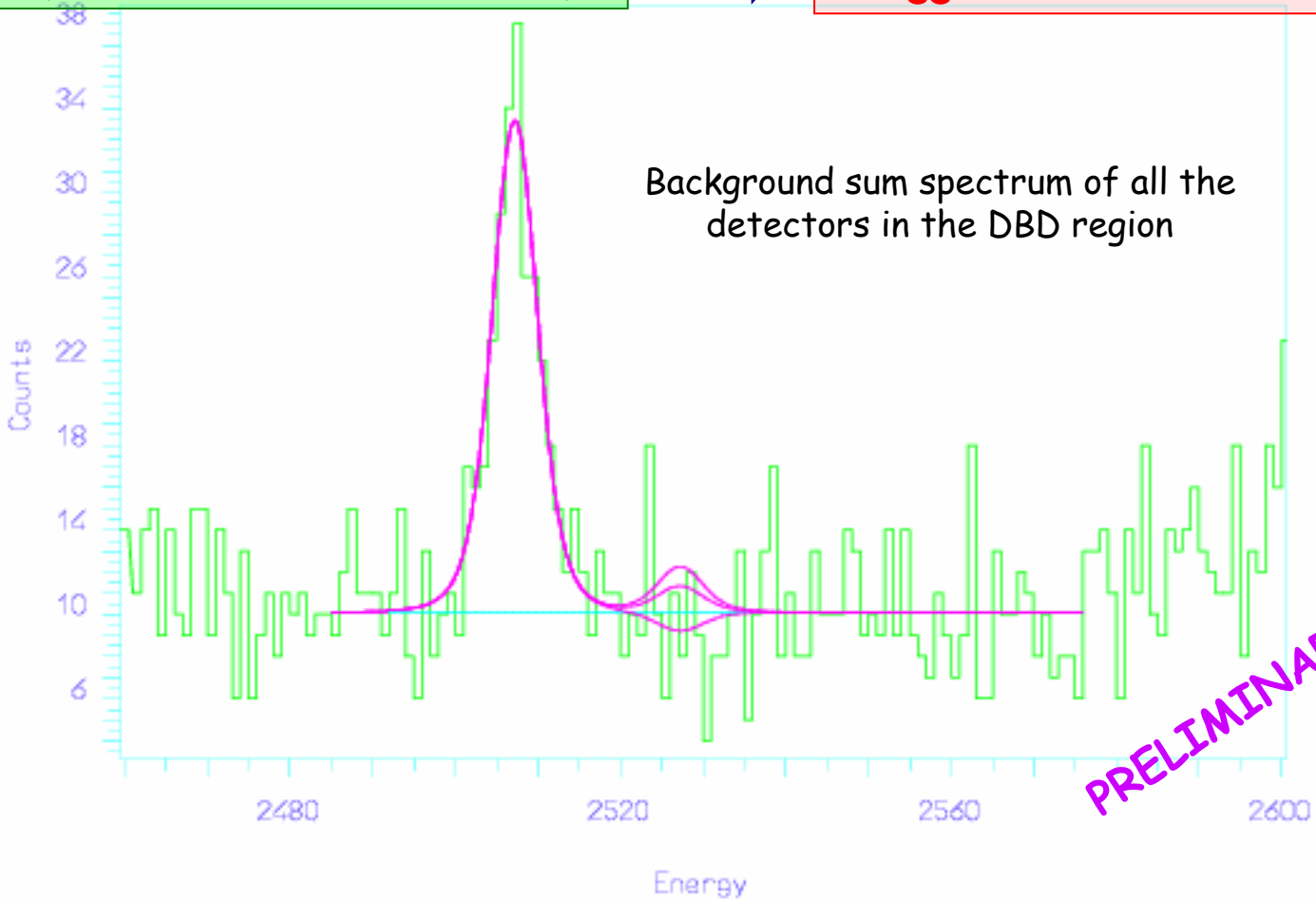
very preliminary

$$MT = 18 \text{ kg } ^{130}\text{Te} \times y$$

$$T_{1/2}^{0\nu} (y) > 2.94 \times 10^{24} y$$

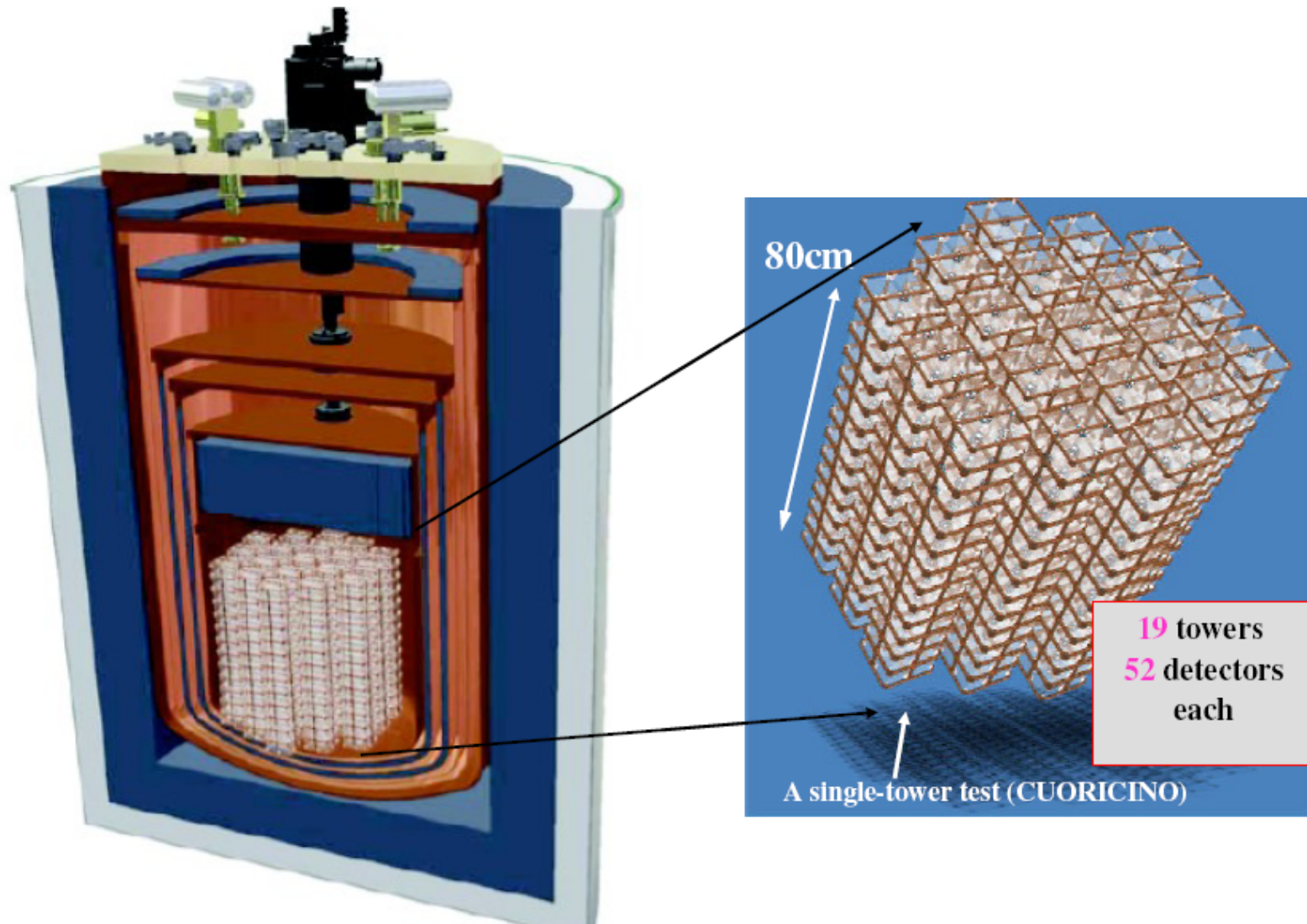


$$m_{ee} < 0.2 - 0.68 \text{ eV}$$



CUORE

Cryogenic Underground Observatory for Rare Events



Closed packed array of 988 TeO_2 $5 \times 5 \times 5$ cm^3 crystals \Rightarrow 741 kg TeO_2 \Rightarrow 204kg ^{130}Te

CUORE status

- CUORE has a **dedicated** infrastructure
- The CUORE **refrigeration** system is the largest in the world
- **1000 crystals** are being installed
- The **first CUORE** crystals were installed in 2010
- CUORE sensitivity $\langle M_{\beta\beta} \rangle \sim 50$ meV with a **factor 2 better** with the new hierarchy region



on has started

en ordered

ch in Nov 2008

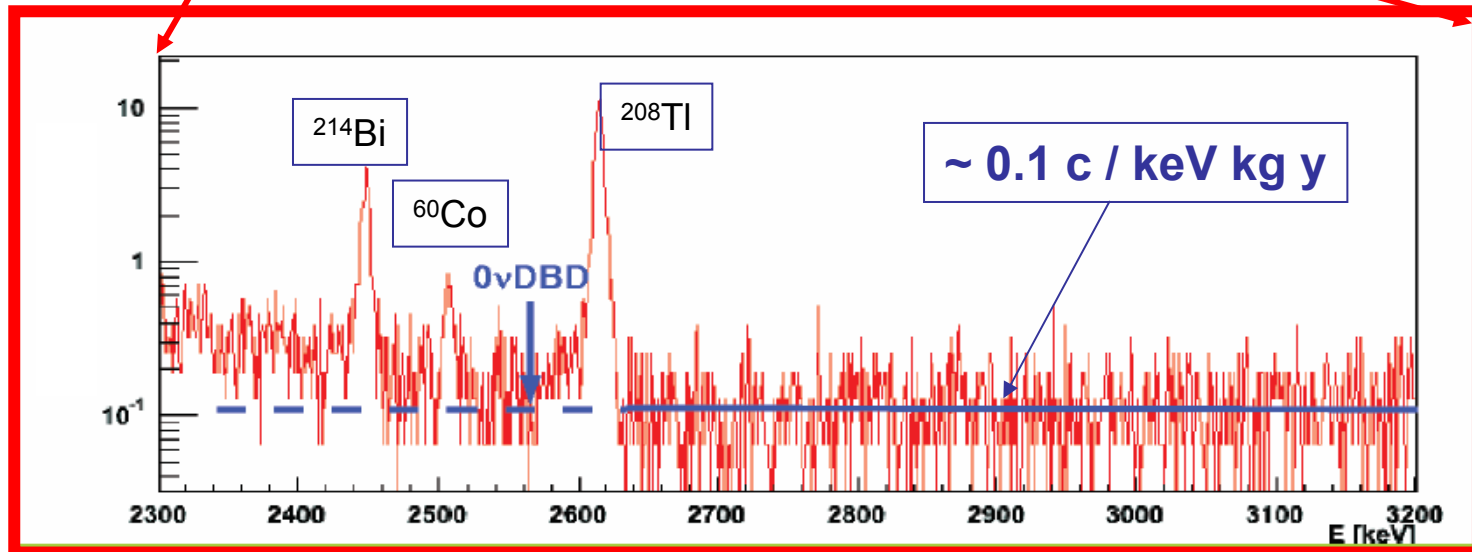
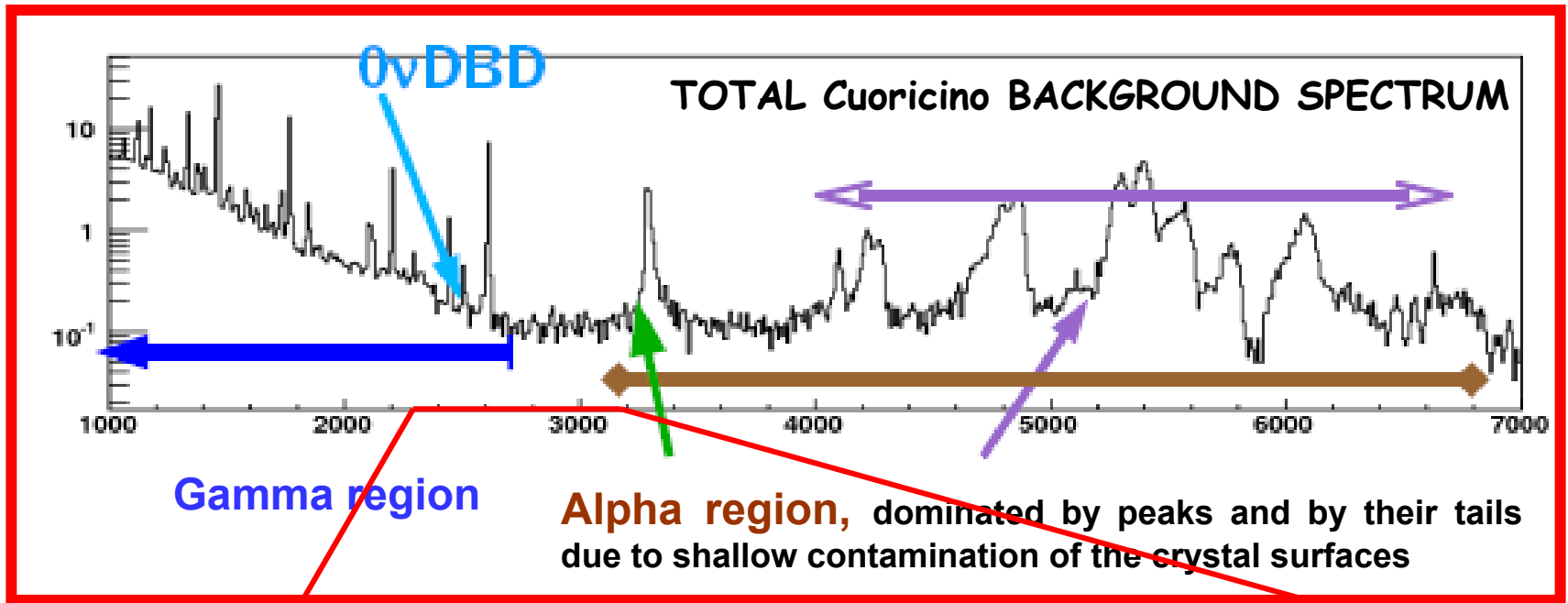
09 and operated

ierarchy region
ted hierarchy region

CUO

2012

Surface background in Cuoricino



Strategies for the control of the surface background from inert materials

(A) Passive methods \Rightarrow surface cleaning

➤ Mechanical cleaning

Crucial test in progress
in the former CUORICINO location

➤ Passivation

(B) Active methods ("reserve weapons" and diagnostic) \Rightarrow events ID

① surface sensitive bolometers

Crucial importance of macrobolometers

→ many interesting high-Q-value candidates
can be studied with very low background

SINGLE BETA DECAY

Model independent tool: the kinematics of β decay

basic idea to measure ν mass: use only kinematics



$$E^2 = M^2c^4 + p^2c^2$$

processes involving neutrinos in the final state

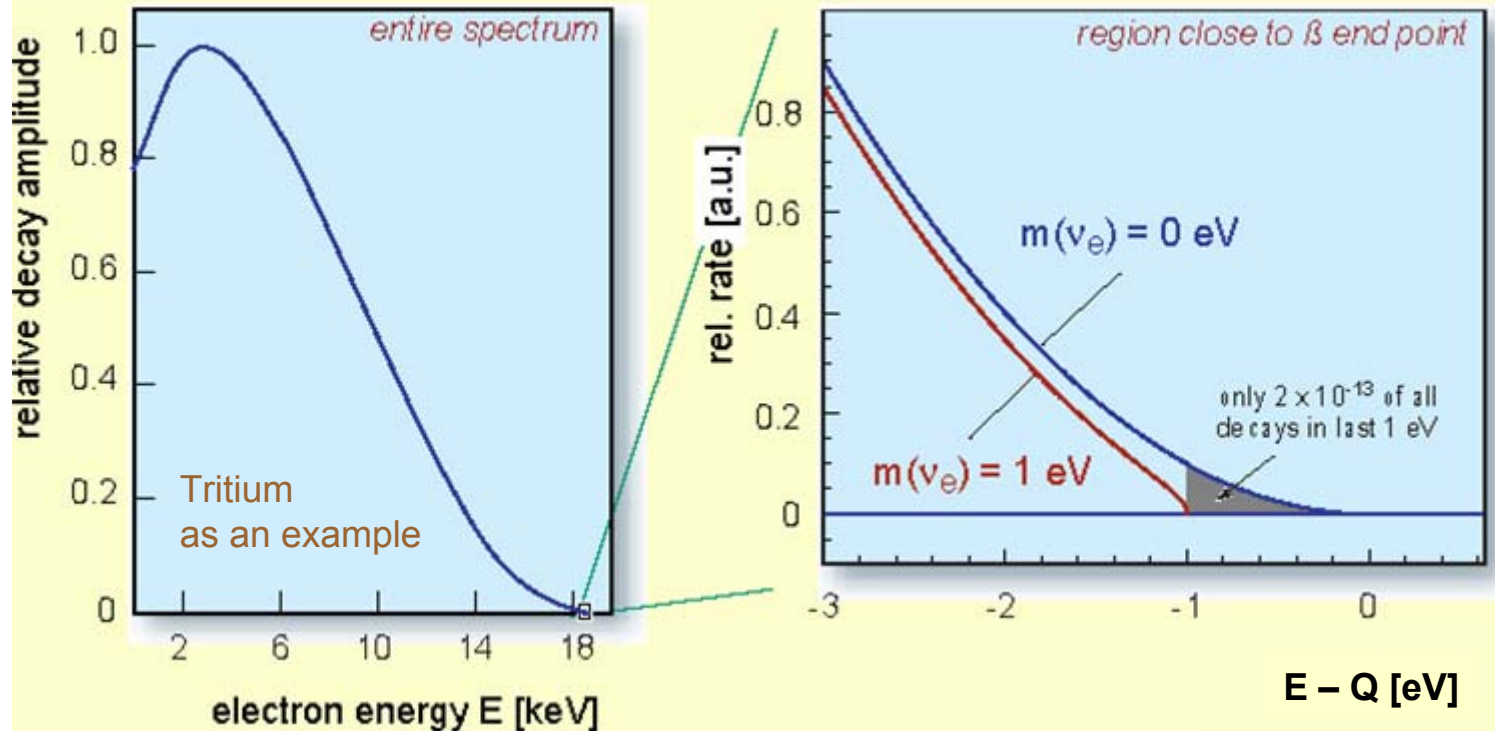
Single β Decay



$$Q = M_{\text{at}}(A,Z) - M_{\text{at}}(A,Z+1) \cong E_e + E_\nu$$

The modified part of the beta spectrum is over range of the order of $[Q - M_\nu c^2, Q]$

Effects of a finite neutrino mass on the beta decay



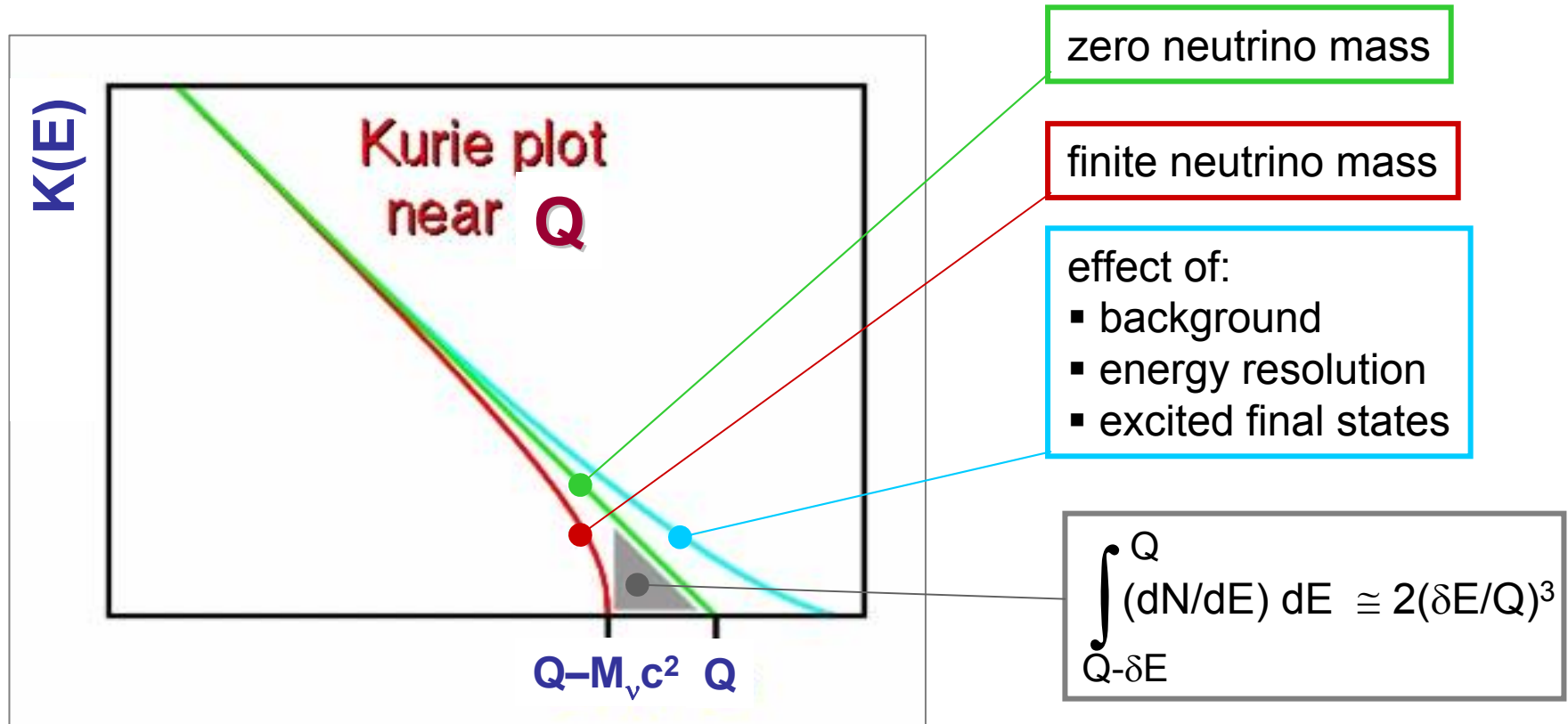
The count fraction laying in this range is $\propto (M_\nu/Q)^3$



low Q are preferred

Effects of a finite neutrino mass on the Kurie plot

The Kurie plot $K(E_e)$ is a convenient **linearization** of the beta spectrum



Mass hierarchy

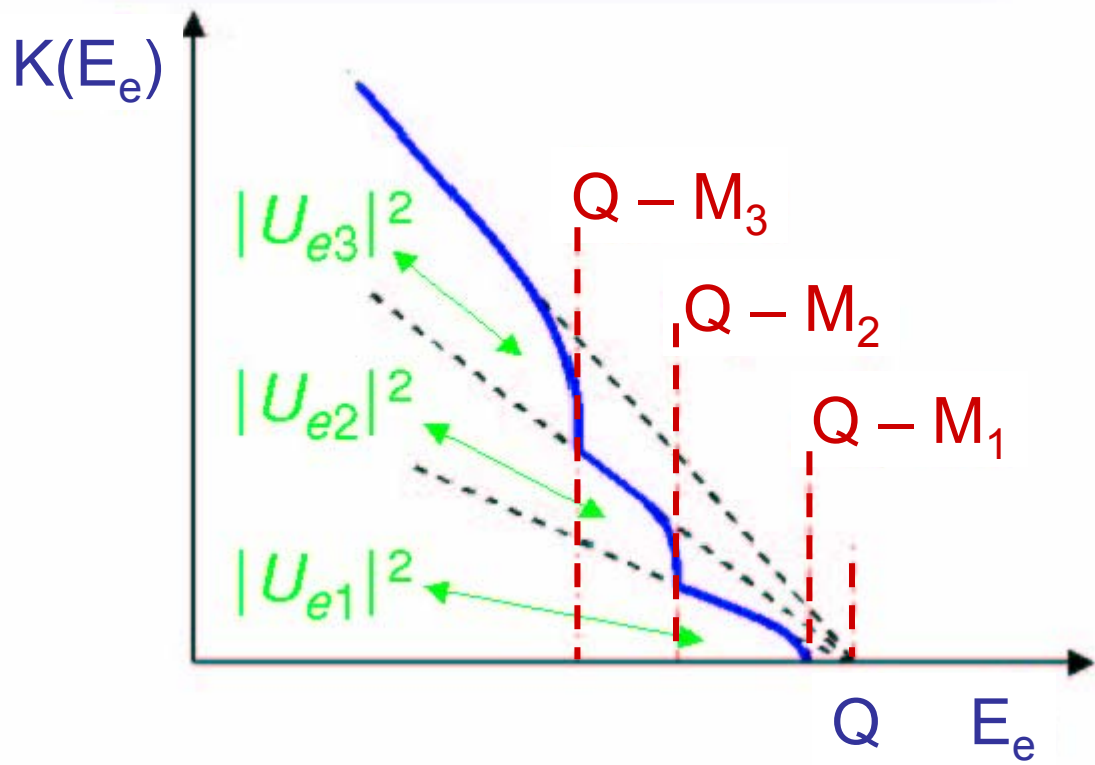
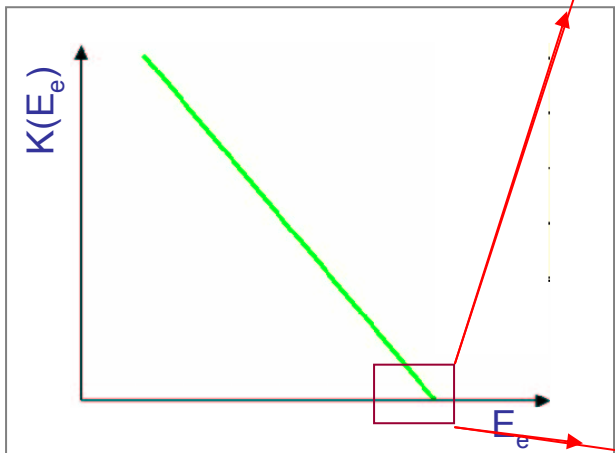
In case of **mass hierarchy**:

- the Kurie plot \equiv superposition of **three different sub - Kurie plots**
- each sub - Kurie plot corresponds to one of the **three different mass eigenvalues**

The weight of each sub - Kurie plot will be given by $|U_{ej}|^2$, where

$$|v_e\rangle = \sum_{i=1}^3 U_{ei} |v_{Mi}\rangle$$

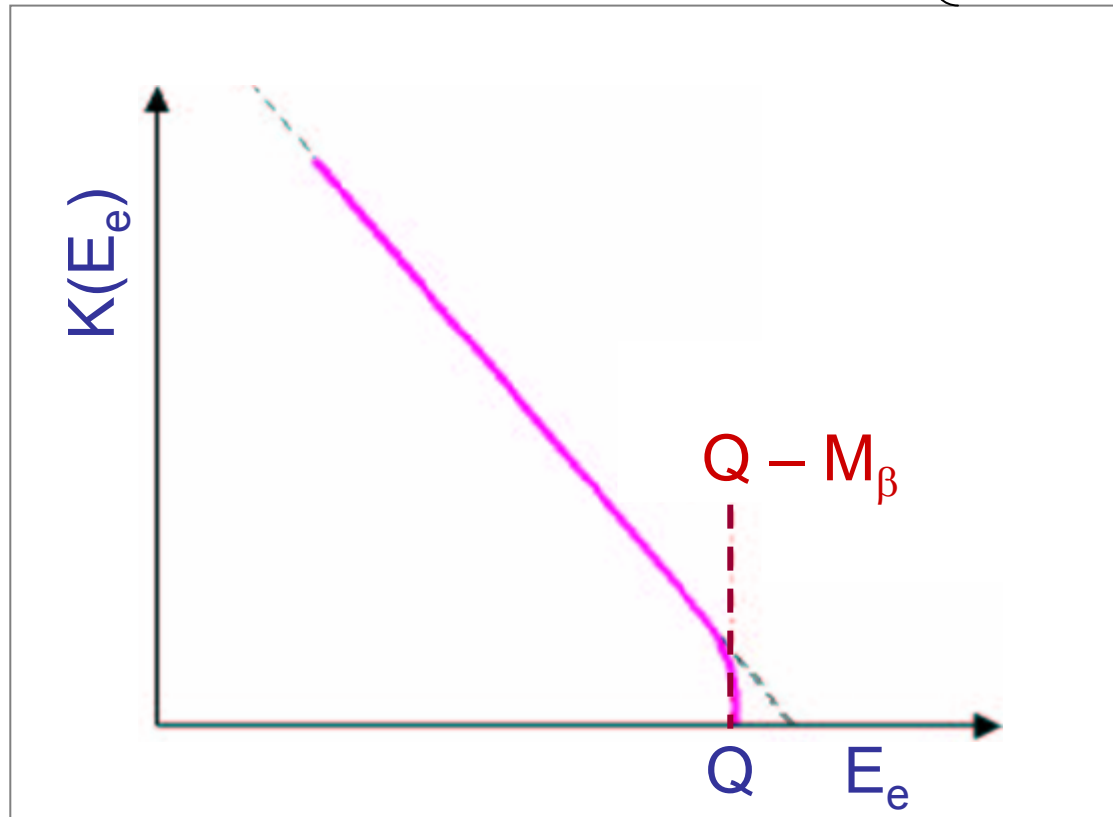
This detailed structure **will not be resolved** with present and planned experimental sensitivities (~ 0.2 eV)



Mass degeneracy

If the 3 mass components **cannot be resolved** or **degeneracy holds**:
the Kurie-plot can be described in terms of a **single mass parameter**,
a mean value of the three mass eigenstates

$$\hookrightarrow \langle M_\beta \rangle = \left(\sum M_i^2 |U_{ei}|^2 \right)^{1/2}$$



Experimental searches based on nuclear beta decay

Requests:

- **high energy resolution** \Rightarrow a tiny spectral distortion must be observed
- **high statistics** in a very narrow region of the beta spectrum
- well known **response of the detector** \Rightarrow spectral output for an energy δ function input
- control of any **systematic effect** that could distort the spectral shape

Approximate approach to evaluate sensitivity to neutrino mass $\sigma(M_\nu)$:

Require that the **deficit of counts** close to the end point due to neutrino mass be equal to the **Poissonian fluctuation** of number of counts in the massless spectrum

The diagram shows the equation for the sensitivity $\sigma(M_\nu)$ to neutrino mass, enclosed in a yellow box. A large blue arrow points from the left towards the equation. Three smaller blue arrows point from labels below to the variables in the equation: 'total source activity' points to 'A', 'live time' points to 'T_M', and 'energy resolution' points to 'ΔE'.

$$\sigma(M_\nu) \cong \sqrt[4]{\frac{1.6 Q^3 \Delta E}{A T_M}}$$

total source activity

live time

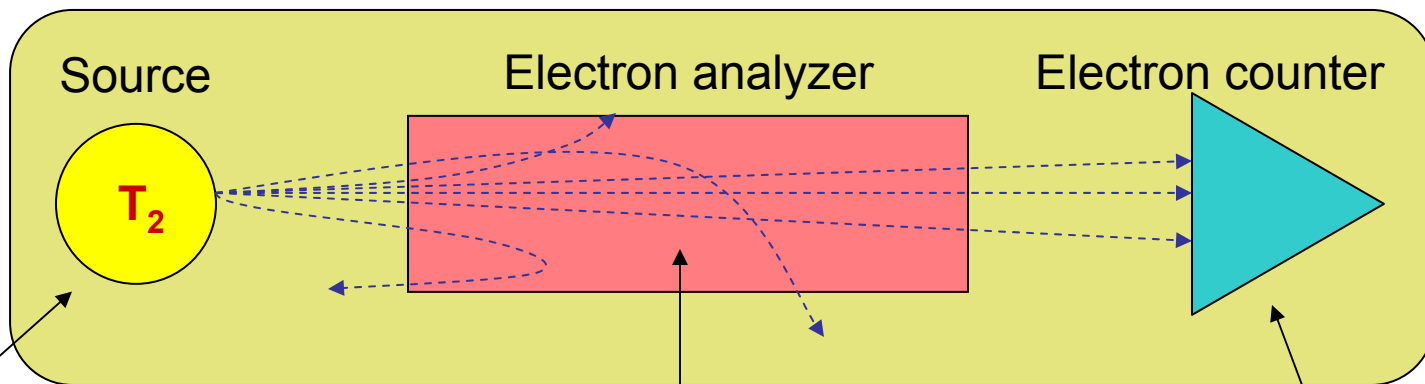
energy resolution

Spectrometers

source **separate** from detector (the source is T - $Q=18.6$ keV)



- determine electron energy by means of a selection on the beta electrons operated by proper **electric and magnetic fields**
- measurement of the electron energy out of the source
- magnetic and electrostatic spectrometers
- present achieved sensitivity: ~ 2 eV (Troitsk/Mainz)
- future planned sensitivity: ~ 0.2 eV (KATRIN)



high activity

high energy resolution
integral spectrum: select $E_e > E_{th}$

- high efficiency
- low background

Microcalorimeters

source \equiv detector (calorimetric approach) (the source is ^{187}Re - $Q=2.5$ keV)

- determine all the "visible" energy of the decay with a high resolution low energy "nuclear" detector
- measurement of the neutrino mass
- cryogenic microcalorimeters
- present achievements
- future prospects

Spectrometers and microcalorimeters have completely different systematic uncertainties

electron
excitation energies

Contra: All the beta electrons are collected
Not-resolved low-energy pulse pile-up
is the main problem
→ it distorts the end-point region

presence of decays to
states, the calorimeter
measures both the electron and
the de-excitation energy

on
dN/dE

Calorimeter requirements

Requirements for a **sensitive calorimetric measurement**:

- precise determination of the β energy
- high statistics
- low pile-up fraction

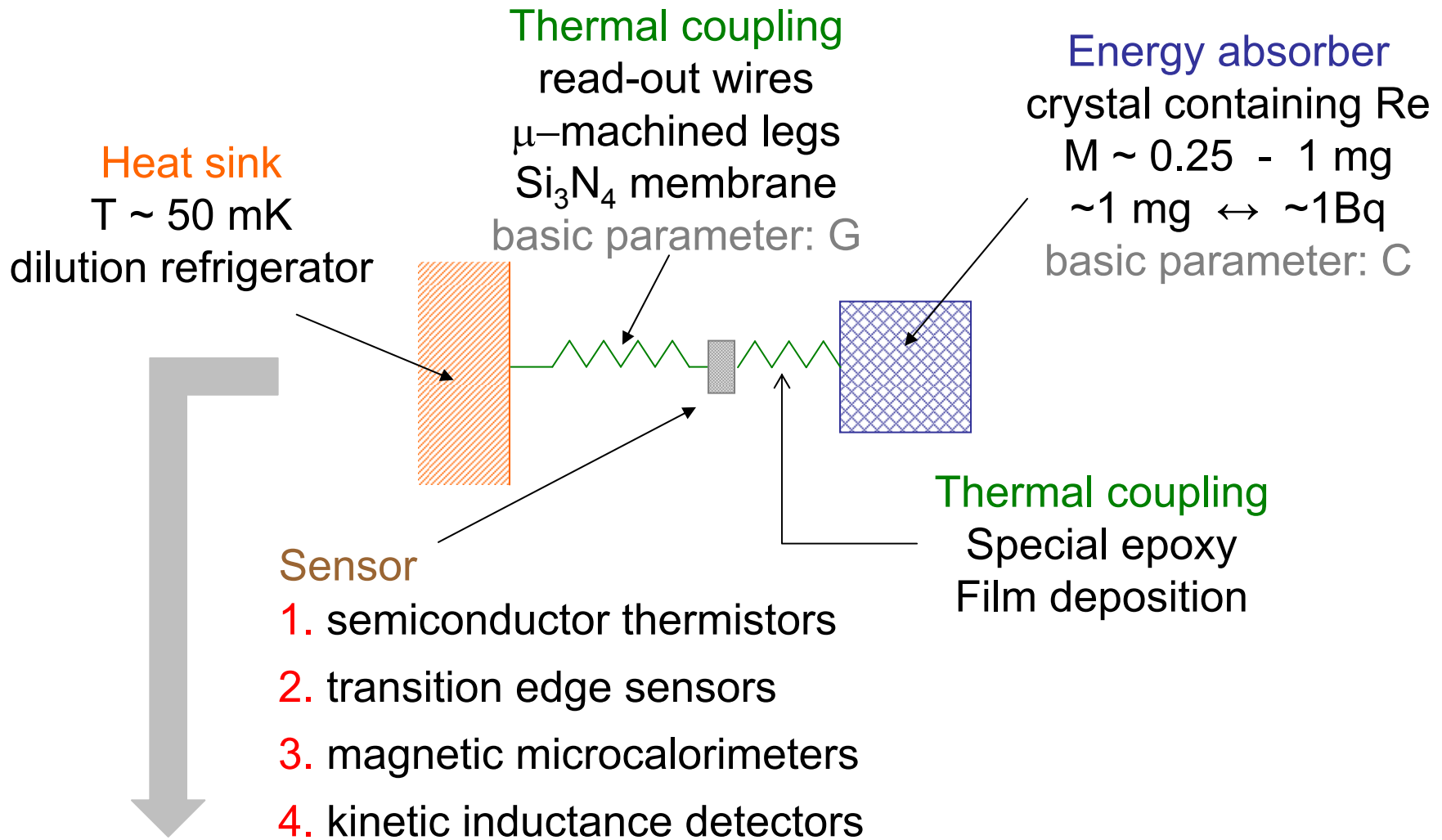


- ✓ short pulse-pair resolving time
- ✓ fractionate the whole detector in many independent elements

In terms of **detector technology**:
development of a single element with these features:

- extremely high energy resolution in the keV range (1‰)
- very fast risetime (present 100 μ s \rightarrow planned 1 μ s)
- high reproducibility of the single element
- possibility of multiplexing

Scheme of a ^{187}Re microcalorimeter



3 - 4 eV energy resolution with optimized absorbers
10-20 eV energy resolution obtained with Re-based absorber

MIBETA (Milano/Como) experiment: the detectors

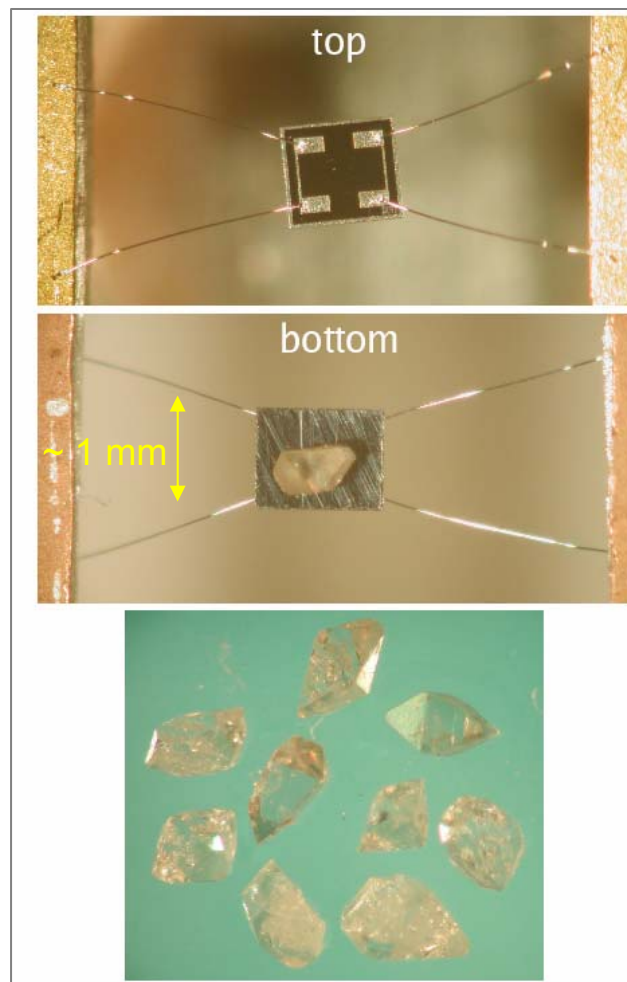
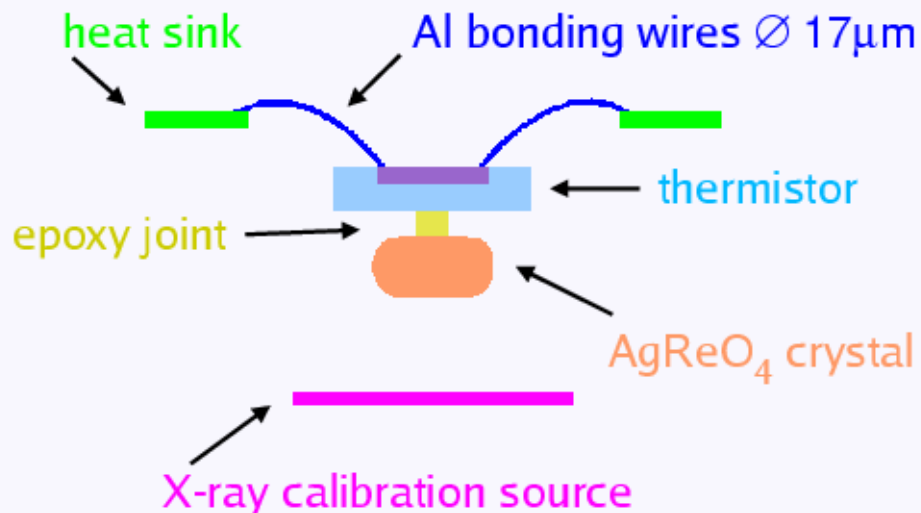
Energy absorbers

- AgReO_4 single crystals
- ^{187}Re activity $\cong 0.54$ Hz/mg
- $M \cong 0.25$ mg $\Rightarrow A \cong 0.13$ Hz

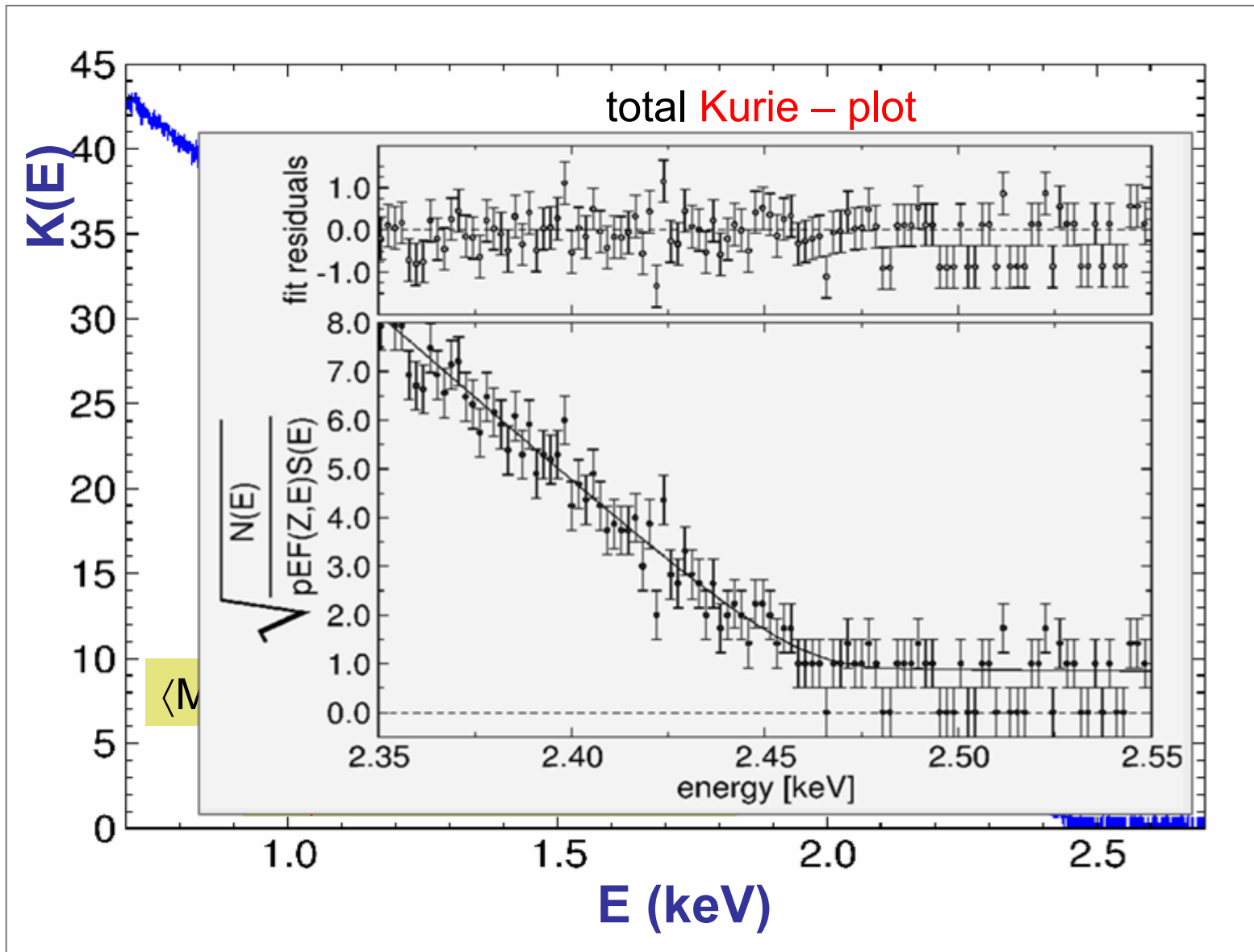
Thermistors

- Si-implanted thermistors
- high sensitivity
- many parameters to play with
- high reproducibility \Rightarrow array
- possibility of μ -machining

typically, array of 10 detectors
lower pile up & higher statistics



MIBETA experiment: the Kurie - plot



The future of bolometric experiments: MARE (Microcalorimeter Arrays for a Rhenium Experiment)

General strategy: push up bolometric technology aiming at:

- multiplication of number of channels
- improvement of energy resolution
- decrease of pulse-pair resolving time

Precursors (MANU, MIBETA)

Semiconductors

Single element
Array of 10 elements
Statistics: $N = 10^6$ events

$\sigma(\langle M_\beta \rangle)$
~
20 eV

MARE-1 - starting - $\Delta E \sim 10 - 30$ eV - $\tau_R \sim 100$ μ s

Transition Edge Sensors
Semiconductors

Arrays of 300 elements
Statistics: $N = 10^{10}$ events

$\sigma(\langle M_\beta \rangle)$
~
2 eV

MARE-2 - 4 years for R&D - $\Delta E \sim 5 - 10$ eV - $\tau_R \sim 1 - 10$ μ s

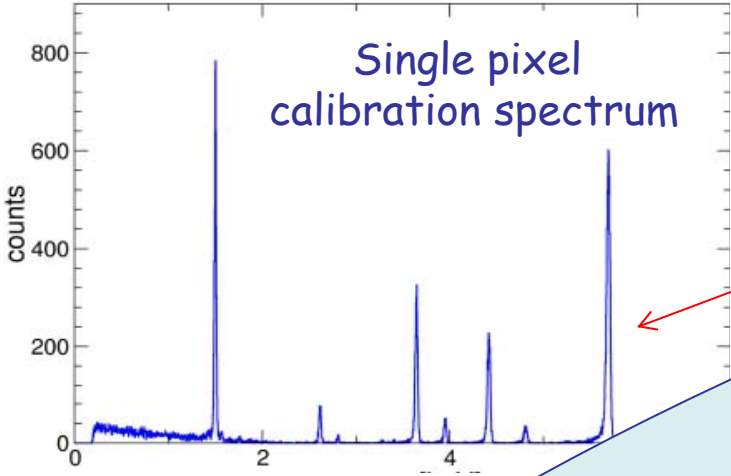
Transition Edge Sensors
Magnetic calorimeters
Kinetic Inductance Det.

Arrays of 50000 elements
Statistics: $N = 10^{13}$ events
Fast τ_R to control pile-up

$\sigma(\langle M_\beta \rangle)$
~
0.2 eV

MARE-1 with Si thermistors: imminent data taking

Single pixel calibration spectrum

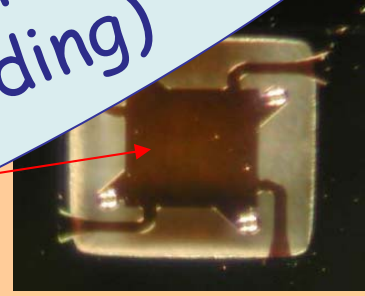
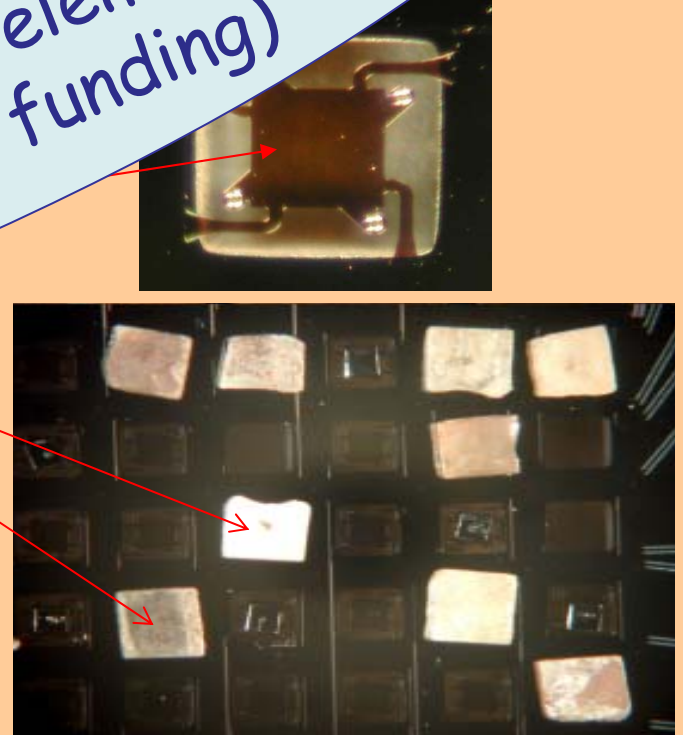


- NASA 6x6 arrays of Si implanted thermistors
- AgReO_4 crystals - 0.4 - 0.8
- Best operation
- $\Delta E \sim 20$

Set-up with 72 elements completed
Data taking before summer
Expandable to 288 elements
(depends only on funding)



AgReO_4 crystals



Conclusions (1)

Why are LT calorimeters so important for neutrino mass and physics?



Two basic features of these devices allow experimental approaches impossible or very difficult with other technologies.

1

Wide choice of the detector materials



Inclusion of the relevant isotopes

2

Extremely high energy resolution

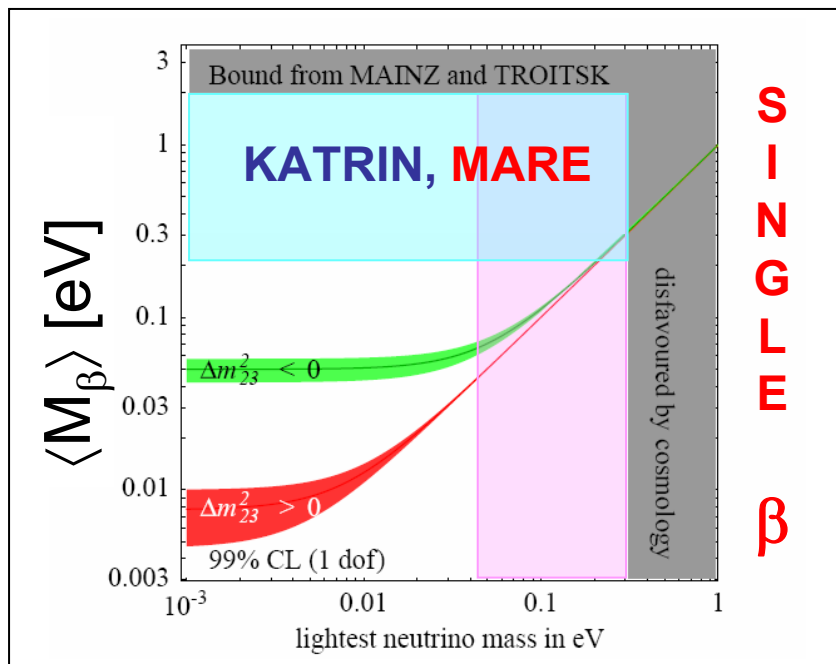


Identification of characteristic spectral features

Conclusions (2)

Exciting times for neutrino masses:

- *degeneracy* will be deeply probed
- discovery potential in case of *inverted hierarchy*



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